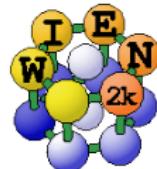


# Optical Properties with Wien2k

Elias Assmann

Vienna University of Technology,  
Institute for Solid State Physics

WIEN2013@PSU, Aug 13



# Menu

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Calculating  $\epsilon$ : Random-Phase Approximation

## ② Practical Calculations

optic: Momentum Matrix Elements

joint: Imaginary Part of Dielectric Tensor

kram: Derived Quantities

## ③ Examples

Ambrosch-Draxl and Sofo, Comp. Phys. Commun. 175, 1 (2006)

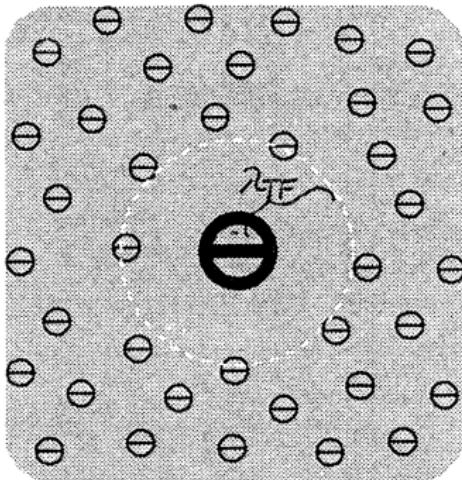
# Appetizer

$$\left\{ \begin{array}{ll} \text{optical conductivity} & \text{Re } \sigma_{ij} = \frac{\omega}{4\pi} \text{Im } \varepsilon_{ij} \\ \text{refractive index} & n_{ii} = \sqrt{(|\varepsilon_{ii}| + \text{Re } \varepsilon_{ii})/2} \\ \text{extinction coefficient} & k_{ii} = \sqrt{(|\varepsilon_{ii}| - \text{Re } \varepsilon_{ii})/2} \\ \text{absorption coefficient} & \alpha_{ii} = 2\omega k/c \\ \text{energy loss function} & L_{ij} = -\text{Im}(\varepsilon^{-1})_{ij} \\ \text{reflectivity} & R_{ii} = \frac{(n-1)^2 + k^2}{(n+1)^2 + k^2} \\ \text{sum rules} & N_{\text{eff}} = \int_0^\omega d\omega' \text{Im } \varepsilon(\omega') \end{array} \right.$$

# Screening

Consider a test charge  $Q$  in a solid:

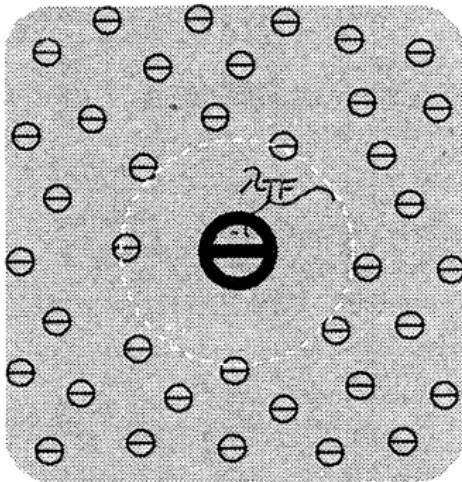
$$V(\mathbf{r} - \mathbf{r}') = \frac{-Q}{|\mathbf{r} - \mathbf{r}'|} \longleftrightarrow V(\mathbf{q}) = -\frac{4\pi Q}{\mathbf{q}^2}$$



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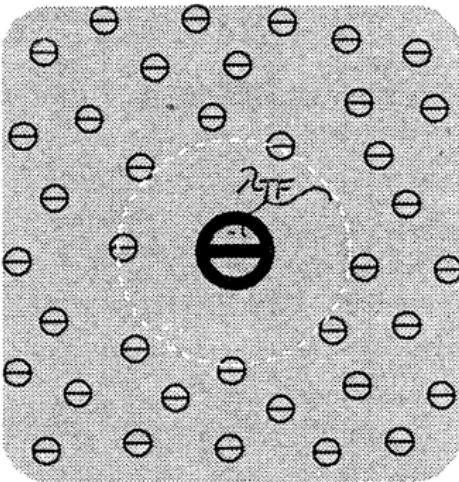


$e^-$  will move to **screen** the charge  
 $\rightsquigarrow$  effective potential  $W$ ;  
dielectric function “ $V = \epsilon W$ ”

# Screening

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→ effective potential  $W$ ;  
dielectric function "  $V = \epsilon W$ "

Simplest model: **Thomas-Fermi**

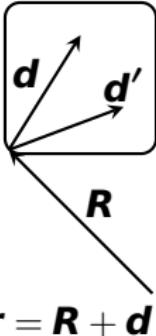
$$W(\mathbf{r}) = \frac{e^{-\mathbf{k}_{\text{TF}}\mathbf{r}}}{\mathbf{r}} \longleftrightarrow W(\mathbf{q}) = \frac{4\pi}{\mathbf{k}_{\text{TF}}^2 + \mathbf{q}^2}$$

$$\mathbf{k}_{\text{TF}}^2 = 4\pi N(E_F)$$

# Ansatz for $\textcolor{blue}{W}$

$$\textcolor{blue}{W} \sim \int d\mathbf{r}' dt \varepsilon^{-1}(\mathbf{r}'; t) \textcolor{orange}{V}(\mathbf{r} - \mathbf{r}'; t - t')$$

# Ansatz for $\mathcal{W}$



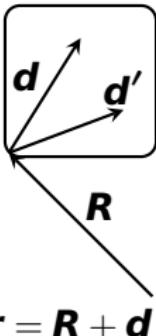
$$\mathcal{W} \sim \int d\mathbf{r}' dt \varepsilon^{-1}(\mathbf{r}'; t) \mathcal{V}(\mathbf{r} - \mathbf{r}'; t - t')$$

Bare  $V(\mathbf{r}, \mathbf{r}'; t, t') = V(\mathbf{r} - \mathbf{r}')\delta(t - t')$  is translation invariant and instantaneous

Response depends on position in unit cell, is retarded

$$\varepsilon^{-1}_{\tilde{\mathbf{R}}}(\mathbf{d}_1, \mathbf{d}_2; t)$$

# Ansatz for $W$



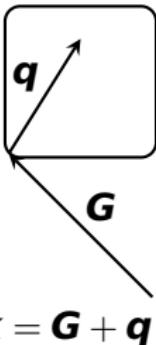
$$W \sim \int d\mathbf{r}' dt \varepsilon^{-1}(\mathbf{r}'; t) V(\mathbf{r} - \mathbf{r}'; t - t')$$

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Response depends on position in unit cell, is retarded

$$\rightsquigarrow W_{\mathbf{R}}(\mathbf{d}, \mathbf{d}'; t) = \sum_{\tilde{\mathbf{R}}} \int d\mathbf{d}_1 d\mathbf{d}_2 \varepsilon^{-1}_{\tilde{\mathbf{R}}}(\mathbf{d}_1, \mathbf{d}_2; t) \\ \cdot V(\mathbf{R} + \mathbf{d} - \mathbf{d}' - [\mathbf{d}_1 - \mathbf{d}_2 - \tilde{\mathbf{R}}])$$

# The Dielectric Function



$$W_G(\mathbf{q}, \omega) = \sum_{\mathbf{G}'} \varepsilon^{-1}_{\mathbf{GG}'}(\mathbf{q}, \omega) V_{\mathbf{G}'}(\mathbf{q}, \omega)$$

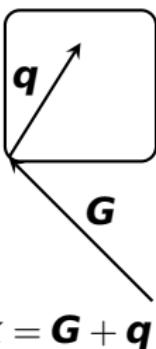
$$\mathbf{k} = \mathbf{G} + \mathbf{q}$$

light is long-wavelength:

$$W_G(\mathbf{q}, \omega) \approx \varepsilon^{-1}_{\mathbf{GO}}(\mathbf{q}, \omega) V_{\mathbf{O}}(\mathbf{q}, \omega)$$

$$\mathbf{G}' = \mathbf{0}, \quad \mathbf{q} \rightarrow \mathbf{0}$$

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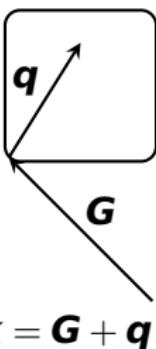
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“macroscopic”  $\varepsilon$  (u.c. average):

$$W(\mathbf{q}, \omega) = \varepsilon^{-1}_{\mathbf{0}\mathbf{0}}(\mathbf{q}, \omega) V_{\mathbf{0}}(\mathbf{q}, \omega)$$

$$\varepsilon_M(\mathbf{q}, \omega) = \frac{1}{\varepsilon^{-1}_{\mathbf{0}\mathbf{0}}(\mathbf{q}, \omega)}$$

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$$\varepsilon_M(\mathbf{q}, \omega) = \frac{1}{\varepsilon^{-1}_{\mathbf{0}\mathbf{0}}(\mathbf{q}, \omega)}$$

neglect local-field effects:

$$\varepsilon_M(\mathbf{q}, \omega) \approx \varepsilon_{\mathbf{0}\mathbf{0}}(\mathbf{q}, \omega)$$

# Calculating $\varepsilon$ : The RPA

$$V(\mathbf{q}) = \varepsilon(\mathbf{q}, \omega) W(\mathbf{q}, \omega)$$

- Poisson:  $\mathbf{q}^2 W = 4\pi(-Q + \delta n) \quad \leftrightarrow \quad W = V + \frac{4\pi}{\mathbf{q}^2} \delta n$
- linear response:  $\delta n = \chi V$

# Calculating $\varepsilon$ : The RPA

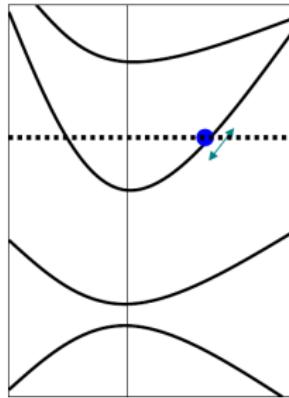
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- Poisson:  $\mathbf{q}^2 W = 4\pi(-Q + \delta n) \leftrightarrow W = V + \frac{4\pi}{\mathbf{q}^2} \delta n$
- linear response:  $\delta n = \cancel{\chi} PW \rightarrow V = (1 - \frac{4\pi}{\mathbf{q}^2} P) W$
- “random-phase” approximation:  $P$  to lowest order

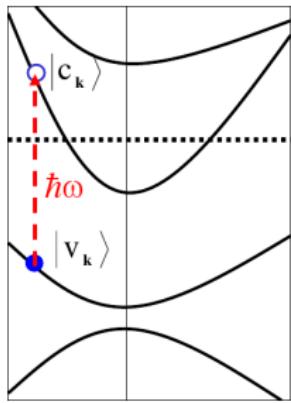
$$P = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots$$

$$\sim G^0(1, 2) G^0(2, 1)$$

# Intra- and Interband transitions



intraband



interband

$E_F$  free  $e^-$ : Lindhard formula

$$P = \frac{4\pi}{q^2\Omega} \sum_{\mathbf{k}} \frac{f(\epsilon_{\mathbf{k}+q}) - f(\epsilon_{\mathbf{k}})}{\epsilon_{\mathbf{k}+q} - \epsilon_{\mathbf{k}} - \omega}$$

Bloch  $e^-$ :

$$P = \frac{4\pi}{q^2\Omega} \sum_{\mathbf{k}nn'} A_{\mathbf{k}\mathbf{q}}^{nn'} \frac{f(\epsilon_{\mathbf{k}+q}) - f(\epsilon_{\mathbf{k}})}{\epsilon_{\mathbf{k}+q} - \epsilon_{\mathbf{k}} - \omega}$$

# Intra- and Interband transitions

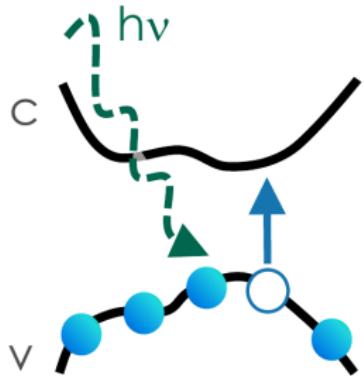
**intraband:** Drude model,  
( $\omega_p$ : plasma frequency)

$$\text{Im } \epsilon^{\text{intra}} = \frac{\Gamma \omega_p^2}{\omega(\omega^2 + \Gamma^2)}$$

**interband:**

joint density of states:

$$\rho(\omega) = \sum_{c,v} \int d\mathbf{k} \delta(\epsilon_c(\mathbf{k}) - \epsilon_v(\mathbf{k}) - \omega)$$

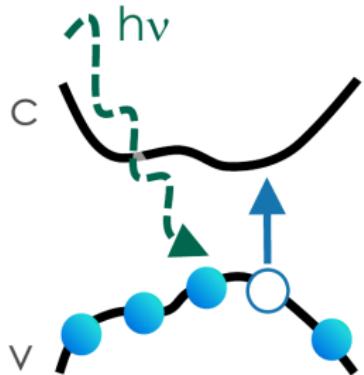


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joint density of states:

$$\rho(\omega) = \sum_{c,v} \int d\mathbf{k} \delta(\epsilon_c(\mathbf{k}) - \epsilon_v(\mathbf{k}) - \omega)$$

$v-c$  transition probability  
("selection rules") given by  
momentum matrix elements

$$\text{Im } \epsilon_{ij}(\omega, \mathbf{0}) \propto \frac{1}{\omega^2} \sum_{c,v} \int d\mathbf{k} \delta(\epsilon_c(\mathbf{k}) - \epsilon_v(\mathbf{k}) - \omega) \langle c\mathbf{k} | \hat{p}_i | v\mathbf{k} \rangle \langle v\mathbf{k} | \hat{p}_j | c\mathbf{k} \rangle$$

# Symmetry Constraints

$\varepsilon_{ij} = \varepsilon_{ji}$  is always symmetric.

Additional constraints from crystal symmetry:

$$\varepsilon = U^{-1} \varepsilon U$$

cubic 
$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

tetragonal,  
trigonal,  
hexagonal 
$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 3 \end{pmatrix}$$

monoclinic 
$$\begin{pmatrix} 1 & 4 & 0 \\ 2 & 0 & 0 \\ 0 & 0 & 3 \end{pmatrix}$$

orthorhombic 
$$\begin{pmatrix} 1 & 0 & 0 \\ 0 & 2 & 0 \\ 0 & 0 & 3 \end{pmatrix}$$

triclinic 
$$\begin{pmatrix} 1 & 4 & 5 \\ 2 & 6 & 0 \\ 0 & 0 & 3 \end{pmatrix}$$

# Program Flow

lapw1 Kohn-Sham eigenstates

optic momentum matrix elements (case.symmat)

joint imaginary part of dielectric tensor (case.joint)

kram derived quantities

- Kramers-Kronig

$$\text{Re } \varepsilon_{ij} = \delta_{ij} + \frac{2}{\pi} \mathcal{P} \int_0^{\infty} d\Omega \frac{\Omega}{\Omega^2 - \omega^2} \text{Im } \varepsilon_{ij}$$

⇒ all optical constants

# optic: Momentum Matrix Elements

- ① normal SCF run → converged density
- ② x `kgen` → dense k-mesh (check convergence!)
- ③ x `lapw1 -options` → eigenvectors on dense mesh
- ④ x `lapw2 -fermi -options` → `case.weight`
  - metals: “TETRA 101.0” in `case.in2`
- ⑤ x `optic -options` → momentum matrix elements  
`case.symmat`:  $\langle c\mathbf{k}|\hat{p}_i|v\mathbf{k}\rangle \langle v\mathbf{k}|\hat{p}_j|c\mathbf{k}\rangle$

core-level spectra: Kevin Jorissen's lecture tomorrow 10:30

# optic: Input and Output

## case.inop

```
99999 1           #k-points, 1st k-point
-5.0 3.0 9999    Emin Emax [Ry], NBvalMAX
2             #indep. elements (symmetry/SOC)
1             Re xx
3             Re zz
OFF 3          write mommat2?, #spheres
1 2 3          spheres to sum over
```

## symmetry

1: Re⟨xx⟩    4: Re⟨xy⟩  
2: Re⟨yy⟩    5: Re⟨xz⟩  
3: Re⟨zz⟩    6: Re⟨yz⟩

## spin-orbit

7: Im⟨xy⟩  
8: Im⟨xz⟩  
9: Im⟨yz⟩

## case.symmat

$\langle v\mathbf{k} | \hat{p}_i | c\mathbf{k} \rangle \langle c\mathbf{k} | \hat{p}_j | v\mathbf{k} \rangle$

## case.mommat2 (if ON)

$\langle v\mathbf{k} | \hat{p}_i | c\mathbf{k} \rangle$

# joint: $\text{Im}(\varepsilon)$ , (Joint) Density of States

## case.injoint

1 9999 9999 lower, upper, upper-val bandindex  
0.0 .001 1.0 Emin ( $\geq 0$ ), dE, Emax [Ry]  
eV units [eV / ryd / cm-1]  
4 mode  
2 #indep. elements  
0.1 0.1 broadenings  $\Gamma$  for Drude (mode=6,7)

## case.joint

$$\left. \rho \right\} \sim \sum_{c,v} \int d\mathbf{k} \delta(\epsilon_c(\mathbf{k}) - \epsilon_v(\mathbf{k}) - \omega) \left\{ \frac{\langle c\mathbf{k} | \hat{p}_i | v\mathbf{k} \rangle \langle v\mathbf{k} | \hat{p}_j | c\mathbf{k} \rangle}{1} \right.$$

# joint: Modes of Operation

“physical” (all bands)

- 1 joint DOS
- 3 regular DOS
- 4  $\text{Im } \epsilon$  interband
- 6  $\text{Im } \epsilon$  intraband (Drude)

$$\text{Im } \epsilon_{ij} \sim \sum_{c,v,\mathbf{k}} \delta(\epsilon_c(\mathbf{k}) - \epsilon_v(\mathbf{k}) - \omega) \langle c\mathbf{k}|\hat{p}_i|v\mathbf{k}\rangle \langle v\mathbf{k}|\hat{p}_j|c\mathbf{k}\rangle$$

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band analysis

- 0 joint DOS
- 2 DOS
- 5 interband
- 7 intraband

$$\text{Im } \epsilon_{ij} \sim \sum_{c,v,\mathbf{k}} \delta(\epsilon_c(\mathbf{k}) - \epsilon_v(\mathbf{k}) - \omega) \langle c\mathbf{k}|\hat{p}_i|v\mathbf{k} \rangle \langle v\mathbf{k}|\hat{p}_j|c\mathbf{k} \rangle$$

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“sphere analysis”

$$|c\mathbf{k}\rangle = \sum_{\alpha}^{\text{MT,I}} |c\mathbf{k}\rangle_{\alpha}$$

NB: cross-terms are missed!

case.inop

```
0FF 3 mommat2?, #spheres  
1 2 3 spheres to sum over
```

# kram: Kramers-Kronig Analysis

## case.inkram

```
0.1      interband broadening
0.0      energy shift (scissors operator)
1        add intraband contributions? 1/0
12.6 12.6 plasma frequencies (joint, mode=6)
0.1  0.1  broadenings  $\Gamma$  for Drude models
```

## output

- `case.epsilon`  $\text{Re } \epsilon, \text{Im } \epsilon$
- `case.sigmak`  $\text{Re } \sigma, \text{Im } \sigma$
- `case.sumrules`
- `case.absorp`  $\text{Re } \alpha, \alpha$
- `case.eloss` loss function
- `case.reflectivity`  $R$
- `case.refraction`  $n, k$

# More Stuff You May Need to Know

## spin-polarized calculations

Kramers-Kronig is not additive.

- ① x joint -up && x joint -dn
- ② adjoint-updn
- ③ x kram

# More Stuff You May Need to Know

## spin-polarized calculations

- ① x joint -up && x joint -dn
- ② addjoint-updn
- ③ x kram

## procedure for metals

- ① x joint (mode=6) → plasma frequencies  $\omega_{pj}$
- ② x joint (mode=4) → interband  $\text{Im } \epsilon$
- ③ x kram (intra=1, insert  $\omega_p$ )

$$\text{Im } \epsilon^{\text{intra}} = \frac{\Gamma \omega_p^2}{\omega(\omega^2 + \Gamma^2)}, \quad \text{Re } \epsilon^{\text{intra}} = 1 - \frac{\omega_p^2}{\omega^2 + \Gamma^2}$$

# More Stuff You May Need to Know

## spin-polarized calculations

- ① x joint -up && x joint -dn
- ② addjoint-updn
- ③ x kram

## procedure for metals

- ① x joint (mode=6) → plasma frequencies  $\omega_{pij}$
- ② x joint (mode=4) → interband  $\text{Im } \epsilon$
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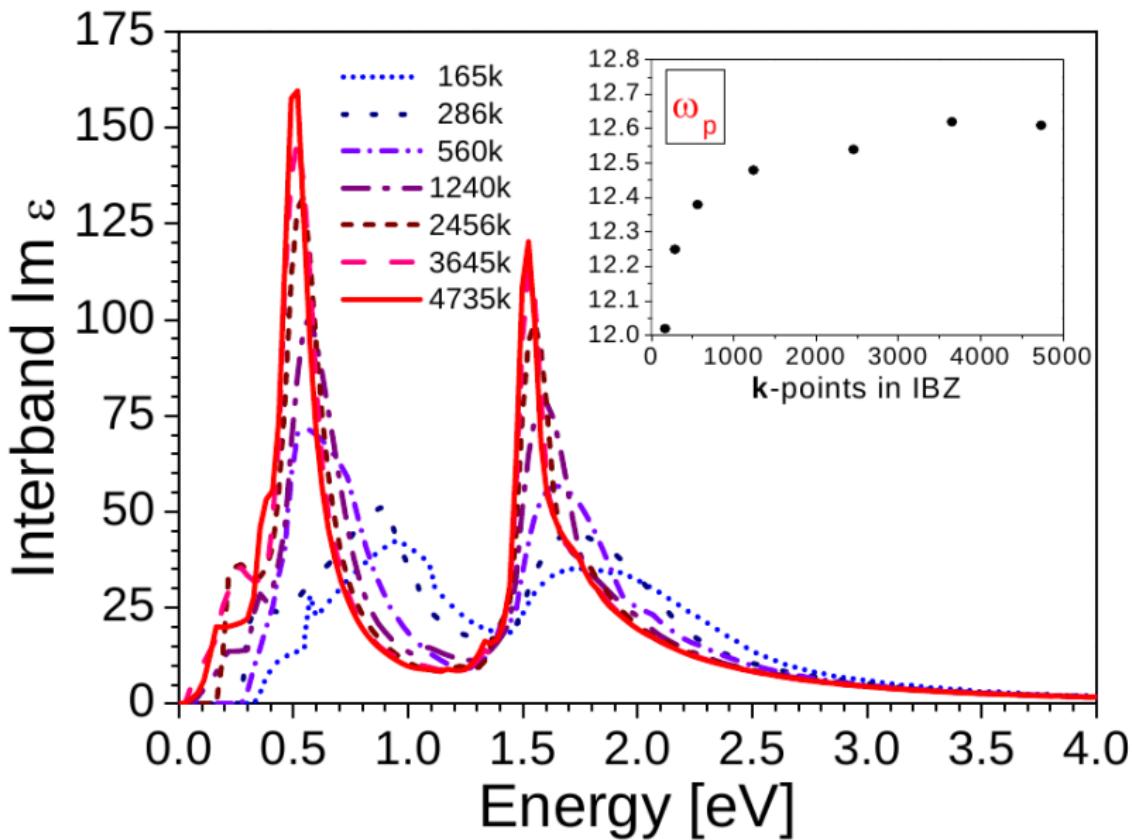
Kramers-Kronig needs  $\text{Im } \epsilon$  in a large energy range

$$\text{Re } \epsilon_{ij} = \delta_{ij} + \frac{2}{\pi} \mathcal{P} \int_0^\infty d\Omega \frac{\Omega}{\Omega^2 - \omega^2} \text{Im } \epsilon_{ij}$$

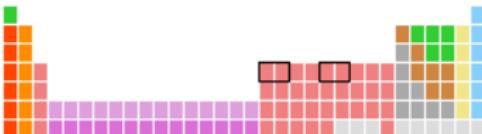
# Some Limitations

- linear optical properties only
  - $W = \varepsilon^{-1(1)} V + \cancel{\varepsilon^{-1(2)} V^2} + \dots$
- Kohn-Sham eigenstates interpreted as excited states
  - ⇒ “scissors” operator:  $\epsilon_c(\mathbf{k}) \rightarrow \epsilon_c^{\text{LDA}}(\mathbf{k}) + \Delta$
- independent-particle approx. (no  $e^- - h^+$  interaction)
  - ⇒ Bethe-Salpeter (BSE) → Peter Blaha’s lecture (13:00)
- LDA/GGA are not exact
  - ⇒ hybrid DFT, effective potentials → Peter Blaha
  - ⇒ DFT+U, LDA+DMFT → my lecture (tomorrow 9:00)

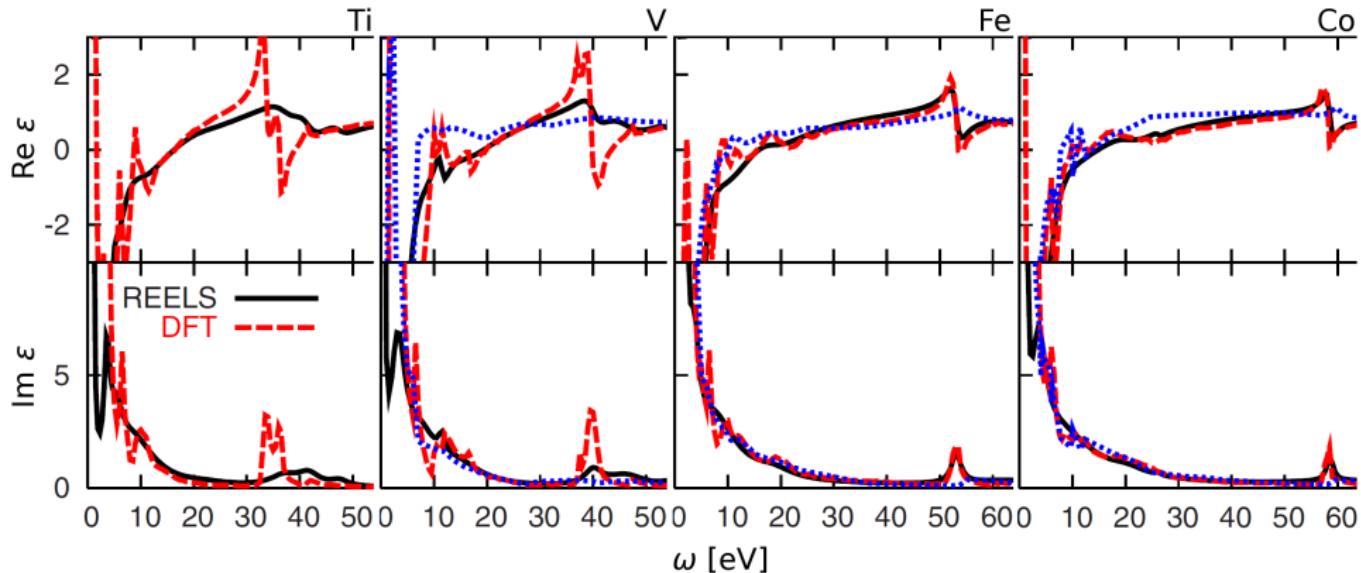
# Example: Al, k-Mesh Convergence



# Comparison to Experiment

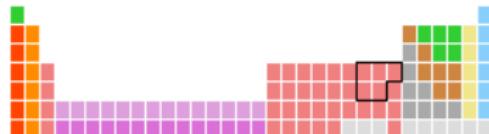


REELS = reflection electron energy loss spectroscopy

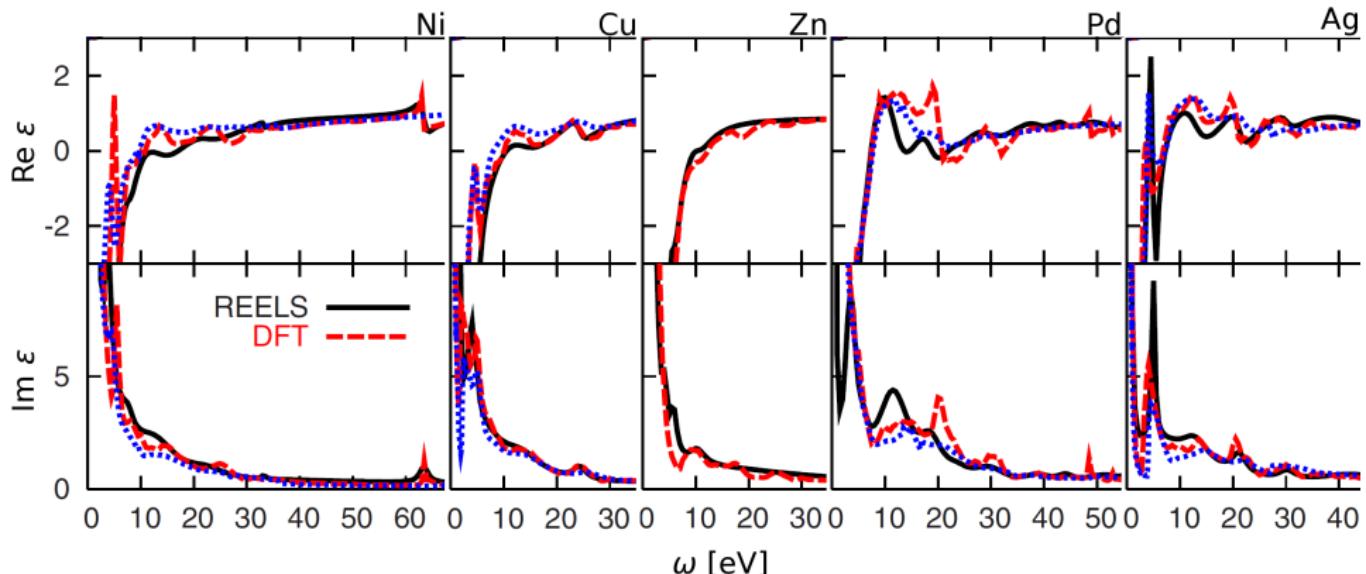


Optical Constants for 17 Elemental Metals  
Werner et al., Phys. Chem. Ref. Data 38, 1013 (2009)

# Comparison to Experiment



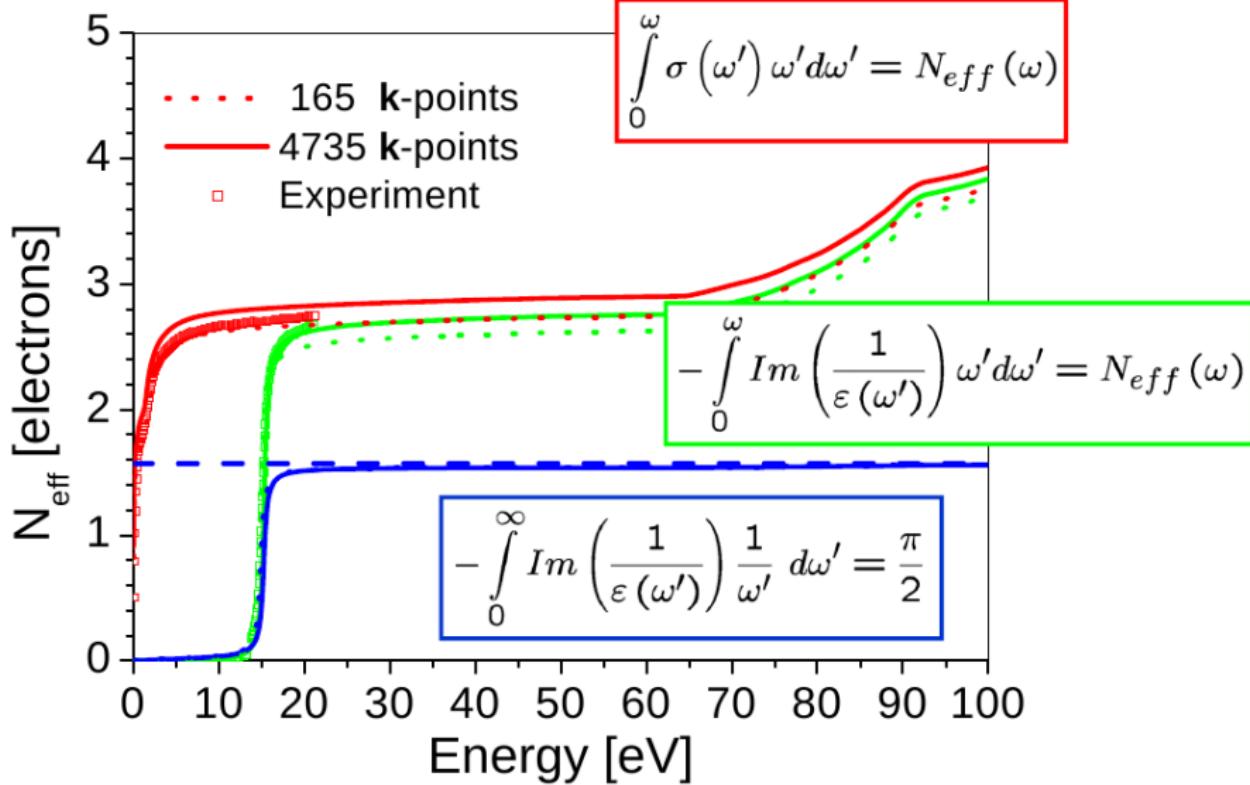
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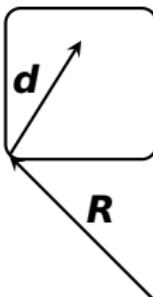
Optical Constants for 17 Elemental Metals

Werner et al., Phys. Chem. Ref. Data 38, 1013 (2009)

# Example: Al, Sum Rules



# Reminder: Lattice Fourier Transform



A function defined on the **unit cell** becomes a discrete function on the **reciprocal lattice**:

$$f(\mathbf{d}) = f(\mathbf{d} + \mathbf{R}) \xleftrightarrow{\mathcal{F}} f_{\mathbf{G}} = \sum_{\mathbf{G}} \delta(\mathbf{k} - \mathbf{G}) f_{\mathbf{G}}$$

$$\mathbf{r} = \mathbf{R} + \mathbf{d}$$
$$\mathbf{k} = \mathbf{G} + \mathbf{q}$$

A discrete function on the **direct lattice** becomes a function defined in the **first BZ**:

$$f_{\mathbf{R}} = \sum_{\mathbf{R}} \delta(\mathbf{r} - \mathbf{R}) f_{\mathbf{R}} \xleftrightarrow{\mathcal{F}} f(\mathbf{q}) = f(\mathbf{q} + \mathbf{G})$$

◀ back