

Stat.Phys.2019: July 3–6, Lviv, Ukraine

Ferromagnetism of LaCoO_3

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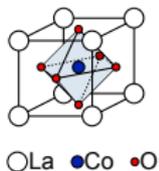
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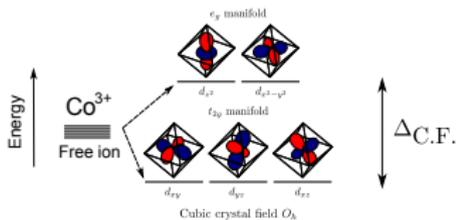
Intro: low- T physics of LaCoO_3 (cubic structure)

The compound is studied since 1960's, both experimentally and theoretically

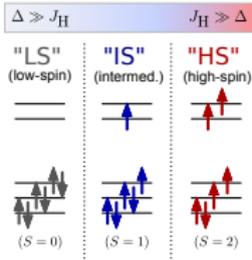
LaCoO_3



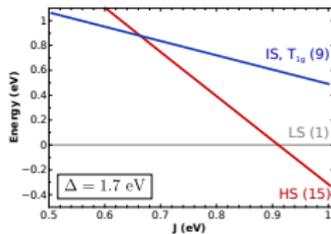
Orbital structure



Relevant spin states



Tanabe-Sugano diagram

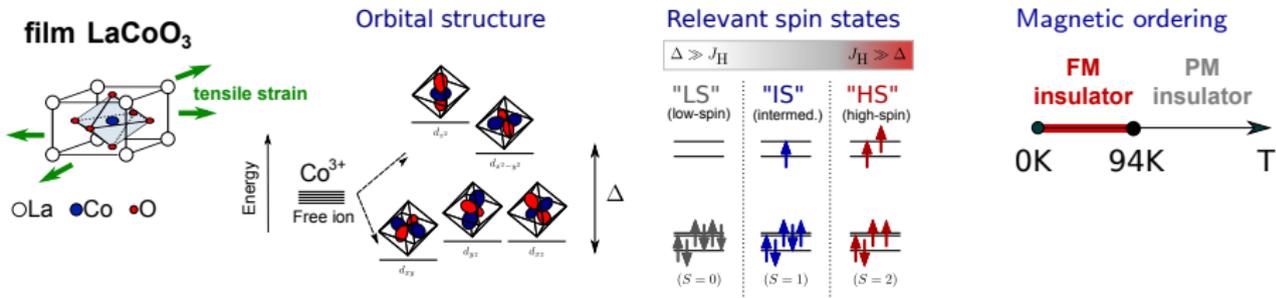


Important facts:

- ✓ at low T : insulator, sizeable charge gap
- ✓ nonmagnetic \rightarrow paramagnetic
- ✓ higher relevance of HS/IS among low- T excitations is still under debate
- ✓ "unusual" behavior at high magnetic fields (~ 100 T)
- ✓ proximity to [A] HS-LS spin-state ordering, [B] BEC of IS ($S = 1$) excitons
- ✓ short-range **ferromagnetic** correlations (no long-range order)



Intro: low- T physics of LaCoO_3 (stretched structure)



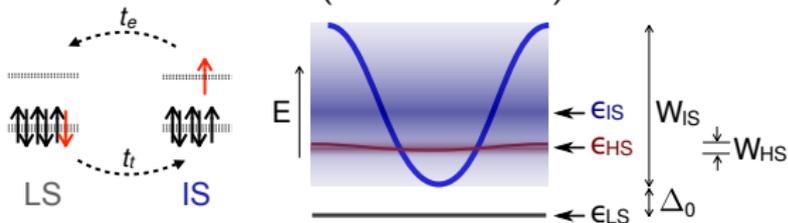
Important facts:

- ✓ insulator
- ✓ long-range ferromagnetic (FM) ordering with $T_C = 94\text{ K}$
- ✓ magnetic structure with a propagation vector $(1/4, -1/4, 1/4)$
- ✓ g factor $\approx 0.7\mu_B$ per Co atom
- ✓ controversial predictions from DFT+U approaches:
 - candidate $\dots\text{HS-LS-LS-LS}\dots$ with magnetic moments on HS sites
 - AFM/FM alignment depends on U ; the FM couplings are too small
- ?? $1/4$ structure, T_C , g factor and the physics behind

Intro: highly-dispersive IS, immobile HS (cubic structure)

LS state \equiv vacuum

IS state \equiv exciton (electron+hole)

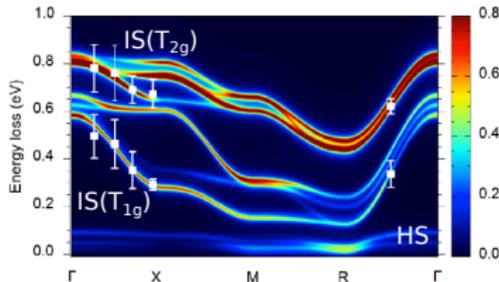
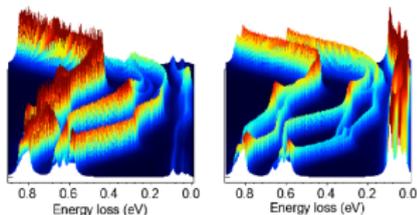


experiment + theory

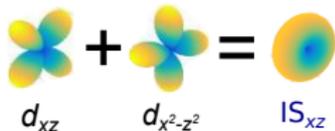
[R.-P. Wang, A. Hariki, *et al.*, PRB **98**, 035149 (2018)]

(a) RIXS

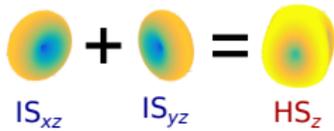
(b) dens. of excit.



IS(T_{1g}) exciton:

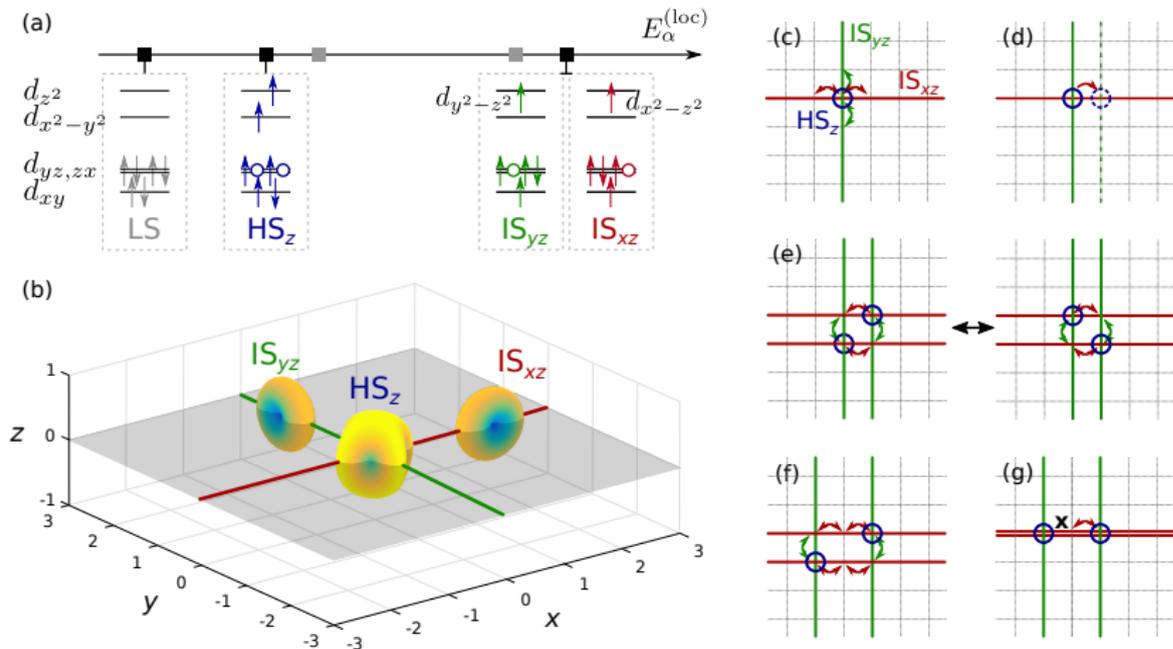


HS(T_{2g}) bi-exciton:



- ✓ $E_{HS} < E_{IS} \Rightarrow$ effective on-site attraction of IS
- ✓ account for formation/decay of HS lowers its energy!

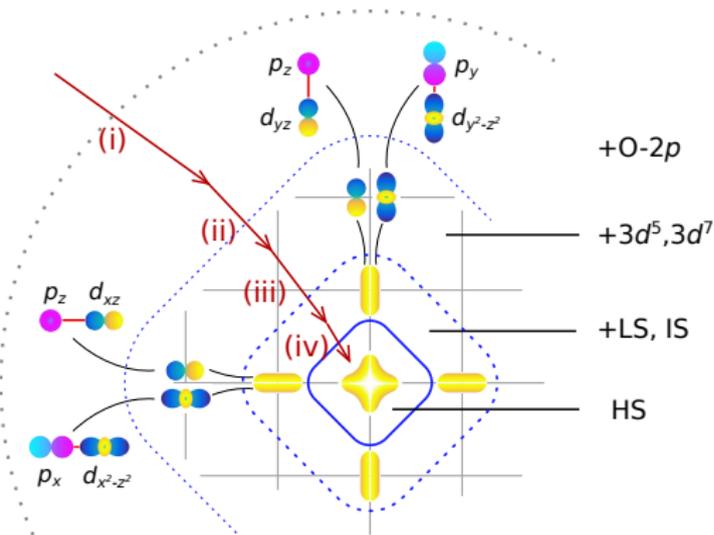
{IS,HS}-boson picture of the film LaCoO_3



Q.: How we extract the hopping and interaction amplitudes for the real crystal?

A.: Series of well-established and logically-consistent approximations...

Methodology



(i) density functional theory (DFT)

[wien2k]

(ii) Wannier projection to d -only Fermi-Hubbard model, adding interaction in the Slater form

[wien2wannier, wannier90]

(iii) strong-coupling (t^2/U) expansion in the d^6 manifold

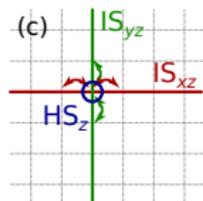
[ED, 2nd order pert. theory]

⇒ effective Bose-Hubbard model
treatment: exact diagonalization

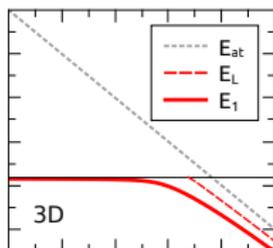
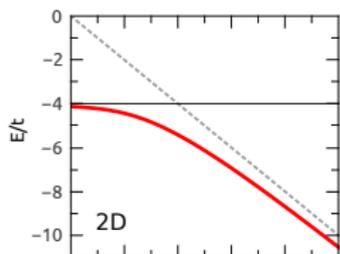
optional, but useful step:

(iv) reduction to the HS-only (spin-1/2) Ising model

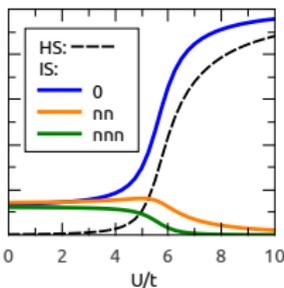
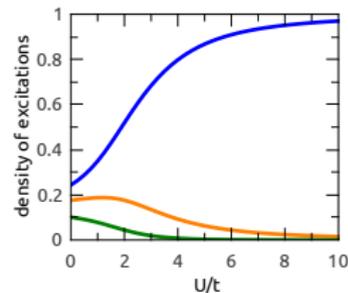
level (iii): “isolated” HS states



- one IS state per line (length ~ 10 sites),
- known IS hopping amplitude t ,
- local attraction U is determined by E_{HS} , E_{IS} (taken as tuneable parameter)

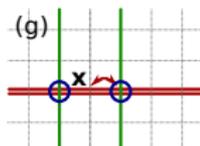


- realistic regime: $U/t \approx 8$
- energy of HS is lowered by nn fluctuations
- HS is surrounded by the IS “cloud” \Rightarrow “dressed” state

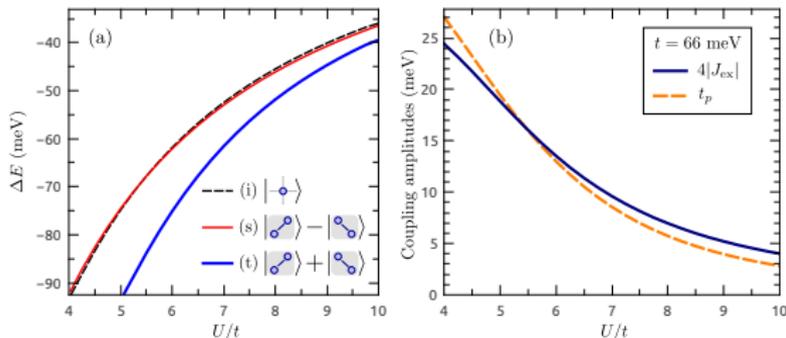
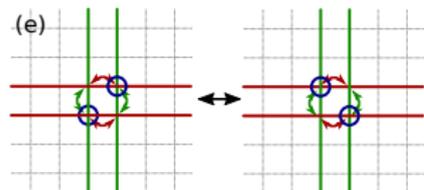


Q.: What if we add excitations?

A.: Rather not on the same line:
(hard-core constraint \rightarrow repulsion)



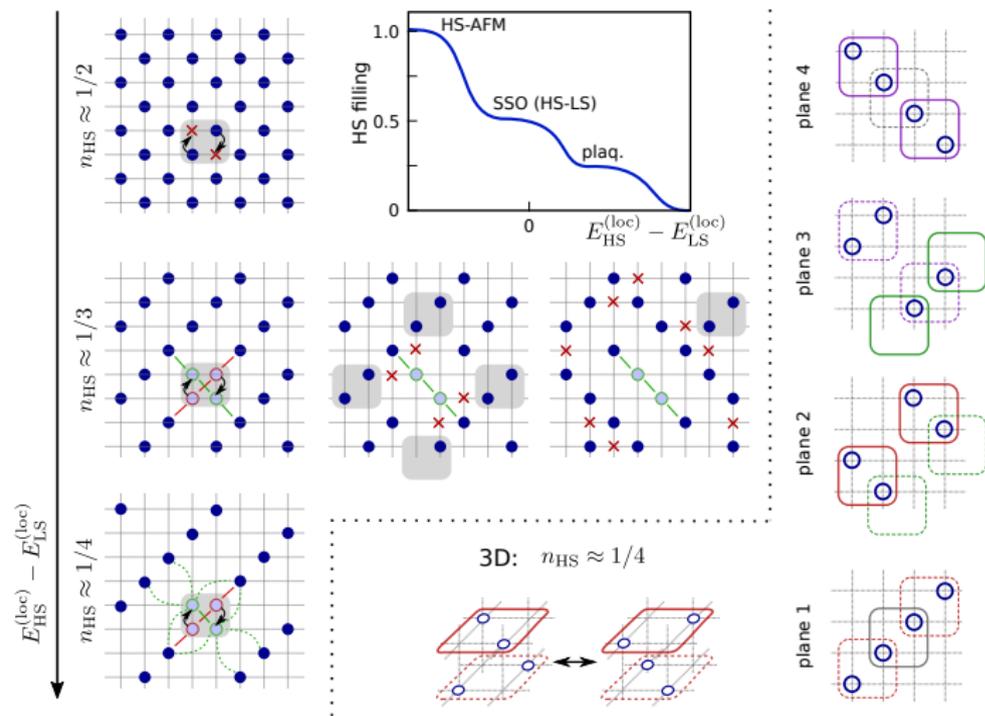
plaquettes with two “dressed” HS & their FM exchange



- effective FM coupling on the diagonal: $4J_{\text{ex}} \approx -7$ meV at $U/t = 8$ ($J_{\text{ex}} < 0$ results from the spin-resolved analysis)
- diagonal $\uparrow\text{-}\uparrow$ attraction $V = 4|J_{\text{ex}}|$ is comparable to $T_c \approx 94$ K
- the amplitude for the pair flip t_p is of the same order

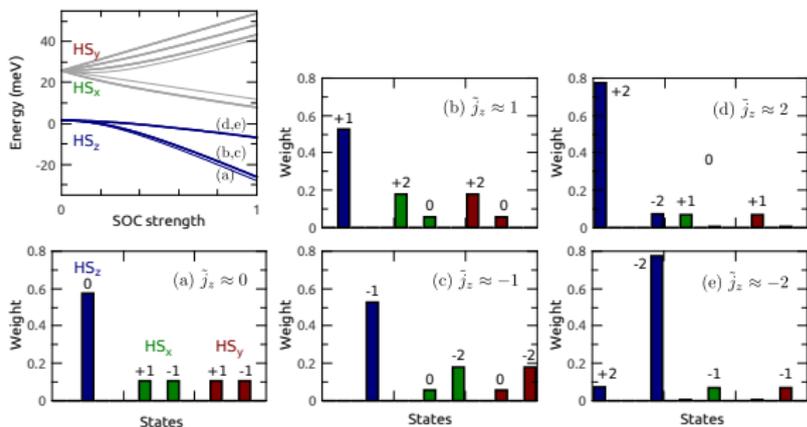
diagonal attraction & nn (nnn) repulsion along $x, y \Rightarrow$ stripe formation!

decorating lattice with the “dressed” HS



supporting remarks: role of SOC; g -factor estimates

Spin-orbit coupling is analyzed in the atomic limit



main effects of SOC:

- filters 3 out of 5 lowest spin states ($\tilde{j} = 0, \pm 1$);
- acts as a single-ion anisotropy term in spin space
- adds a portion of HS_{*x,y*} in the HS_{*z*} set

g factors in the subspace of HS with $\tilde{j} = 0, \pm 1$:

$$M_r = \mu_B(2s_r + l_r) = \mu_B g_r \tilde{j}_r$$

$$g_{\parallel} \approx 2.73$$

$$g_{\perp} \approx 2.64$$

(per HS atom)

Experiment:

$$g_{Co} \approx 0.7 \pm 0.1$$

$$g_{Co} \approx \frac{1}{4} g_{HS}$$

Summary / Acknowledgements

Brief \sum_{slides}

- FM in the insulating film LaCoO_3 originates from the highly-fluctuative character of excitations;
- the lattice is filled by $1/4$ of HS; these form diagonal stripes (agrees with exp.-observed propagation vector)
- FM couplings on diagonals are of the same order as exp. T_C
- at $n_{\text{HS}} = 1/4$, g factors fall close to exp. values
- the low-energy physics of film $\text{LaCoO}_3 \Leftrightarrow$ the (extended) Ising model

Special thanks to:

- | | |
|-----------------------------|--------------------------------|
| ✓ Jan Kuneš (TU Wien); | ✓ Atsushi Hariki (TU Wien); |
| ✓ Kyo-Hoon Ahn (TU Wien); | ✓ Ru-Pan Wang (Utrecht Uni); |
| ✓ Juan F. Afonso (TU Wien); | ✓ Frank de Groot (Utrecht Uni) |

Thank you for your attention!