

Preparation of c-axis fibre textured $\text{YNi}_2\text{B}_2\text{C}$ thin films by pulsed laser deposition: film structure and superconducting properties

K. Häse, S. C. Wimbush, S. Paschen and B. Holzapfel

Abstract—Thin films of $\text{YNi}_2\text{B}_2\text{C}$ have been prepared in situ by pulsed laser deposition onto $\text{MgO}(100)$ substrates. The films exhibit c-axis fibre textured but without in-plane texture. The superconducting transition temperature T_c increases with increasing deposition temperature up to 14.7K at a deposition temperature of 735°C. We observed high values of the upper critical field H_{c2} in thin films of about 7T at 3K and the expected temperature dependence of H_{c2} . An anisotropic behavior of H_{c2} was measured for the different temperatures 2K, 6K and 10K. The critical current density j_c shows different dependencies in applied magnetic fields H_{llc} and $H_{\perp c}$ at 3K.

Index Terms—borocarbides, thin films, upper critical field

I. INTRODUCTION

THE class of quaternary intermetallic superconductors of the $\text{RNi}_2\text{B}_2\text{C}$ (R - rare earth) type, first discovered in 1994 [1, 2], is of primary interest due to its exhibition of the interplay between superconductivity and magnetic ordering phenomena. The series displays several arrangements of magnetic ordering temperature and superconducting transition temperature, including no magnetic ordering, which is the case for $\text{YNi}_2\text{B}_2\text{C}$. This compound is ideal, therefore, to study the fundamental superconducting properties of the series, uninfluenced by magnetic effects. The nonmagnetic members of the series were initially classified as conventional BCS type superconductors having a moderately large density of states at the Fermi level, close to the traditional A-15 superconductors, and a relatively strong electron-phonon coupling constant [3, 4]. In recent years with the availability of high-quality single crystal

samples, the situation has changed somewhat as some properties of $\text{YNi}_2\text{B}_2\text{C}$ and $\text{LuNi}_2\text{B}_2\text{C}$ were observed which might be interpreted as indications of unconventional (*d*-wave) superconductivity [5], for example an anisotropy of the upper critical field within the basal plane of $\text{LuNi}_2\text{B}_2\text{C}$, a quadratic flux line lattice at high magnetic fields and an unconventional behaviour of the specific heat [for an overview see Ref. 6. and references therein]. Several groups are involved in the preparation of thin films of these materials by sputtering and pulsed laser deposition techniques [7, 8, 9]. The availability of c-axis oriented or epitaxial thin films enables important experiments to be performed probing the anisotropic properties of the material. Here, we report on the preparation and anisotropic superconducting properties of c-axis fibre textured $\text{YNi}_2\text{B}_2\text{C}$ thin films prepared by pulsed laser deposition.

II. EXPERIMENTAL

Thin films of $\text{YNi}_2\text{B}_2\text{C}$ were deposited using a KrF excimer laser (wavelength 248nm) in an ultra high vacuum system (base pressure 2×10^{-7} Pa). The deposition conditions comprised a target-substrate separation of 4cm, a laser pulse repetition rate of 10Hz, and an incident energy density of 8 J/cm² onto the target. Polycrystalline $\text{YNi}_2\text{B}_2\text{C}$ targets were obtained by arc melting of a stoichiometric mixture of elemental powders of Y (purity 99.9%), Ni (99.99%), B (99.0%) and C (99.9%) in an inert atmosphere of flowing argon. Single crystal $\text{MgO}(100)$ substrates were positioned

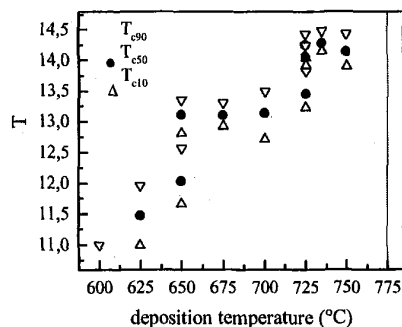


Fig. 1. Inductively measured transition temperature of $\text{YNi}_2\text{B}_2\text{C}$ thin films vs. deposition temperature

Manuscript received September 18, 2000. This work was supported in part by the Deutsche Forschungsgemeinschaft via SFB463.

K. Häse is with the Institute for Solid State and Material Research, PF 270016, D-01171 Dresden, Germany (telephone: +49 351 4549 203, e-mail: k.haese@ifw-dresden.de)

S. C. Wimbush is with the Institute for Solid State and Material Research, PF 270016, D-01171 Dresden, Germany (telephone: +49 351 4549 220, e-mail: s.c.wimbush@ifw-dresden.de)

S. Paschen is with Max Planck Institute for Chemical Physics of Solids, Nöthnitzer Str. 48, D-01187 Dresden, Germany (telephone: +49 351 46463125, email: paschen@cpfs.mpg.de)

B. Holzapfel is with the Institute for Solid State and Material Research, PF 270016, D-01171 Dresden, Germany (telephone: +49 351 4549 455, e-mail: b.holzapfel@ifw-dresden.de)

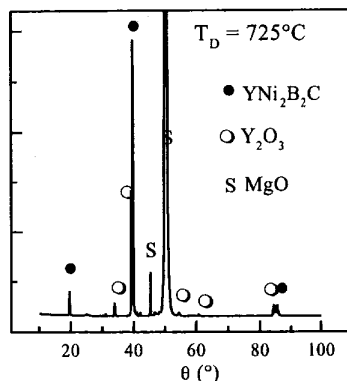


Fig. 2. X-ray diffraction pattern in Bragg-Brentano geometry ($\text{CoK}\alpha$ radiation) of a $\text{YNi}_2\text{B}_2\text{C}$ film deposited at 725°C (● $\text{YNi}_2\text{B}_2\text{C}$, ○ Y_2O_3 and S substrate reflections)

in an on-axis geometry and heated with a radiation heater. The deposition temperature was estimated using a thermocouple positioned at the rear of the heater. The crystal structure of the films was determined using X-ray diffraction in Bragg-Brentano geometry ($\text{CoK}\alpha$ radiation), by grazing incidence diffraction ($\text{CuK}\alpha$) and texture measurements ($\text{CuK}\alpha$). The microstructure was examined with scanning electron microscopy (SEM). Transition temperatures were inductively measured by a magnetic alternating field shielding technique. Transition temperatures measured by this technique are systematically lower than resistively measured transition temperatures. Electrical transport measurements were performed using a four-point contact method on a photolithographically patterned bridge, which was $20\mu\text{m}$ wide and $600\mu\text{m}$ long, in a superconducting 8T split-coil magnet system.

III. RESULTS AND DISCUSSION

Optimisation of the borocarbide phase formation was achieved by varying the deposition temperature. The superconducting transition temperature T_c increases with

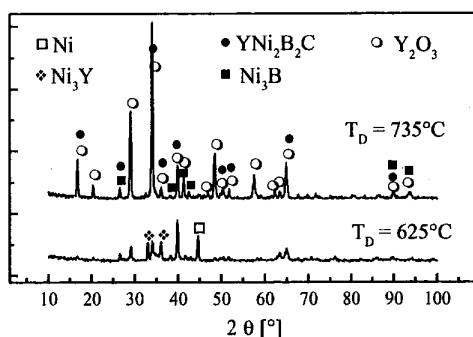


Fig. 3. Grazing incidence x-ray diffraction patterns of $\text{YNi}_2\text{B}_2\text{C}$ films deposited at 625°C and 735°C ($\text{CuK}\alpha$ radiation, incident angle $\omega=4^\circ$)

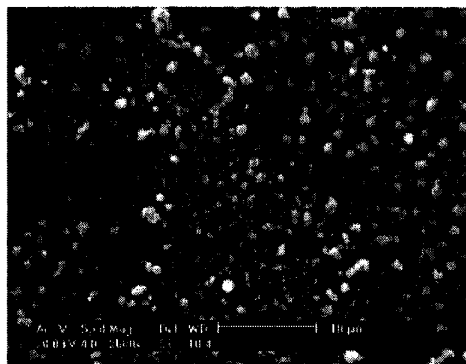


Fig. 4. SEM image of a $\text{YNi}_2\text{B}_2\text{C}$ film deposited at 725°C

increasing deposition temperature (Fig. 1) and the highest superconducting transition temperature was found at 735°C . At higher temperatures, films began to crack and peel away from the substrate upon cooling. X-ray diffraction patterns in Bragg-Brentano geometry (Fig.2) show sharp $(00l)$ borocarbide peaks [10], indicating a strong c-axis orientation, as well as peaks corresponding to secondary phases. We identified the secondary phases Y_2O_3 and Ni_3B by grazing incidence X-ray diffraction from a film deposited at 625°C (Fig.3). At lower temperatures, for instance at 625°C , no borocarbide peaks were observed (Fig.3). In the $\text{YNi}_2\text{B}_2\text{C}$ films the Y_2O_3 phase formation has been shown to take place at the substrate-film interface [11]. X-ray pole figures of the (112) peak did not reveal any in-plane texture for these films. In-plane textured films with comparable superconducting properties have been prepared using lower energy densities ($3\text{J}/\text{cm}^2$), a reduced substrate-target separation and repetition rates of 30Hz.

The microstructure of the films was investigated by SEM. In the $\text{YNi}_2\text{B}_2\text{C}$ films a majority of small round borocarbide grains are visible, as well as a number of larger, irregularly shaped impurity phase regions (Fig. 4). The grain size is about 400nm at a deposition temperature of 725°C (Fig.4).

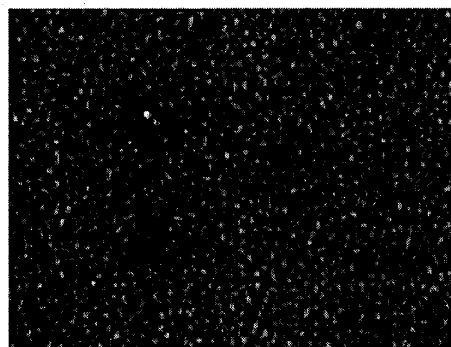


Fig. 5. SEM image of an $\text{YNi}_2\text{B}_2\text{C}$ film deposited at 600°C ; scan area $11.5\mu\text{m} \times 9.5\mu\text{m}$

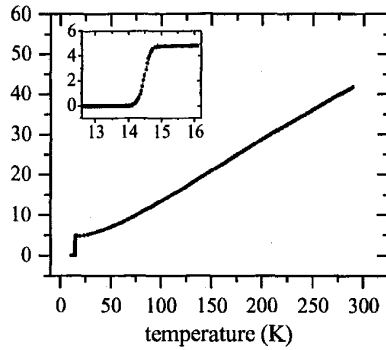


Fig. 6. Temperature dependence of the resistivity for an $\text{YNi}_2\text{B}_2\text{C}$ film deposited at 725°C ; inset: superconducting transition in detail

At lower deposition temperatures, 600°C in Fig. 5, the grains are significantly smaller and less impurity phase regions are visible. Fig. 6 shows the temperature dependence of the resistivity of an $\text{YNi}_2\text{B}_2\text{C}$ thin film. The film shows a superconducting transition at 14.7K with a transition width of 0.3K and a residual resistance ratio ($R_{300\text{K}}/R_{\text{just above } T_c}$) of 9. The decrease in resistivity is linear between 300K and 75K ; in the low temperature range the temperature dependence changes to a quadratic behavior. The resistivity just above T_c is $5\mu\Omega\text{cm}$ (best single crystal values: $1.3\mu\Omega\text{cm}$ [12]).

To determine the temperature dependence of the upper critical field H_{c2} , the $R(T)$ transitions were measured under various magnetic fields applied both parallel and perpendicular to the c -axis. The expected shift in T_c with increasing field is observed as well as a successive broadening of the transition width from 0.3K at 0T to 1.5K at 6T (Fig. 7). The temperature dependence of the upper critical field H_{c2} was obtained from these resistivity curves using the midpoints of the transitions. The general behavior of the temperature dependence of H_{c2} (Fig. 8) is comparable

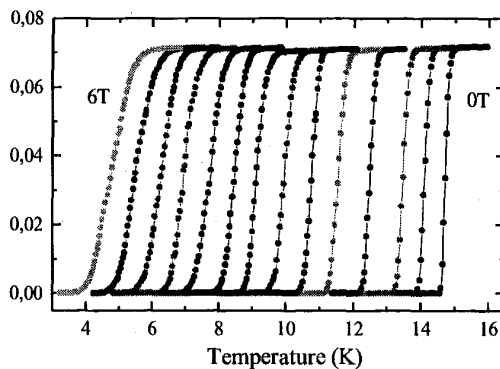


Fig. 7. Temperature dependence of the resistivity for various applied magnetic fields $H_{||c}$ (0, 0.25, 0.5, 1, 1.5, 2, 2.5, 3, 3.5, 4, 4.5, 5, 5.5 and 6T) for an $\text{YNi}_2\text{B}_2\text{C}$ film deposited at 725°C

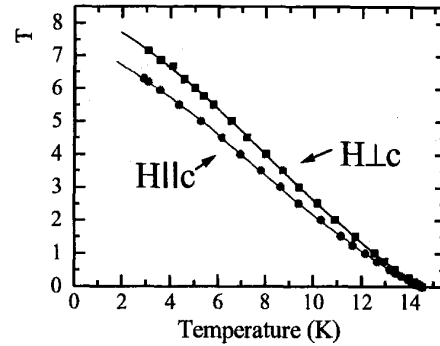


Fig. 8. Temperature dependence of the upper critical field H_{c2} for an $\text{YNi}_2\text{B}_2\text{C}$ film deposited at 725°C

with single crystal data [12] and can be described by the multi-band Eliashberg theory [13]. An anisotropic behavior of H_{c2} was found in $\text{YNi}_2\text{B}_2\text{C}$ for the two different applied field directions (Fig. 8). Measurements of the angular dependence of H_{c2} were carried out to determine the nature of the anisotropic behavior. The magnetic field dependence of the resistivity was measured at temperatures of 2K , 6K and 10K for various angles between the c -axis and the applied field ($0^\circ\text{--}180^\circ$). The angular dependence of H_{c2} is shown in Fig. 9 and its general form was found to be independent of temperature, showing a slight broadening of the curve as the

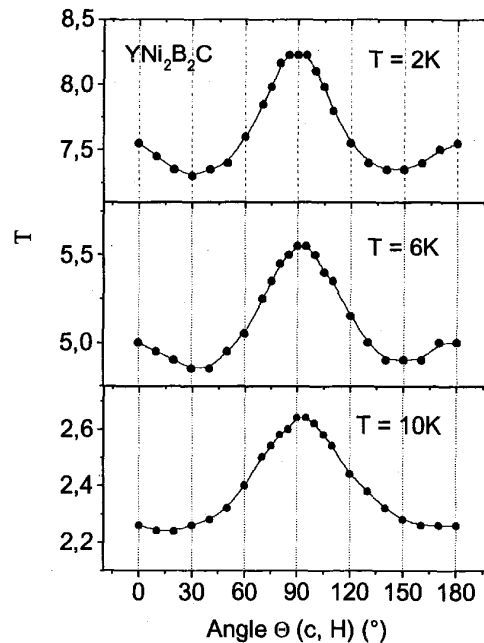


Fig. 9. Angular dependence of the upper critical field H_{c2} for an $\text{YNi}_2\text{B}_2\text{C}$ film at various temperatures (θ - angle between c -axis and applied field)

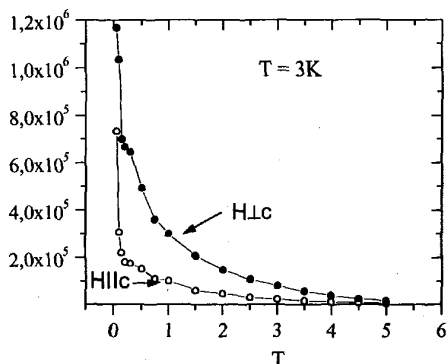


Fig. 10. Magnetic field dependence of the critical current density for an $\text{YNi}_2\text{B}_2\text{C}$ film at 3K for the magnetic field parallel and perpendicular to the c -axis

temperature increased. There is a clear temperature independent minimum at $20^\circ\text{-}30^\circ$. The complex and anisotropic Fermi surface of $\text{RNi}_2\text{B}_2\text{C}$ calculated by Singh [14] and Drechsler *et al.* [15], can account for this unconventional angular dependence of H_{c2} .

The critical current density j_c was measured at a fixed temperature of 3K under an applied magnetic field (Fig. 10). In zero field a high critical current density of $1.2 \times 10^6 \text{ A/cm}^2$ was found although j_c decreases severely in the increasing magnetic field. Both the upper critical field and the critical current density show higher values if the magnetic field is applied perpendicular to the c -axis of $\text{YNi}_2\text{B}_2\text{C}$.

ACKNOWLEDGMENT

The authors would like to thank G. Behr for the preparation of the polycrystalline targets. We wish to thank especially S.-L. Drechsler and H. Rosner for many helpful discussions about our experimental results.

REFERENCES

- [1] Nagarajan, R., Mazumdar, C., Hossain, Z., Dhar, S.K., Gopalakrishnan, K.V., Gupta, L.C., Godart, C., Padalia, B.D. and Vijayaraghavan, Phys. Rev. Lett. 72, 274 (1994)
- [2] Cava, R.J., Takagi, H., Zandbergen, H.W., Krajewski, J.J., Peck Jr., W.F., Siegrist, T., Batlogg, B., van Dover, R.B., Felder, R.J., Mizuhashi, K., Lee, J.O., Eisaki, H., and Uchida, S., Nature 367, 254 (1994)
- [3] L. F. Mattheiss, Phys. Rev. B 49, 13279 (1994)
- [4] W. E. Pickett and D. J. Singh, Phys. Rev. Lett. 72, 3702 (1994)
- [5] G. Wang and K. Maki, Phys. Rev. B 58, 6493 (1998)
- [6] S.-L. Drechsler *et al.*, Physica C 317-318, 117 (1999)
- [7] Häse, K., Holzapfel, B. and Schultz, L., Physica C 288, 28-32 (1997)
- [8] Arisawa, S., Hatano, T., Hirata, K., Mochinu, T., Kitaguchi, H., Fujii, H., Kumakura, H., Kadowaki, K., Nakamura, K. and Togano, K., Appl. Phys. Lett. 65, 1299 (1994),
- [9] Andreone, A., Iavarone, M., Vaglio, R., Manini, P. and Cogliati, E., Appl. Phys. Lett. 69, 118 (1996)
- [10] Siegrist, T., Zandbergen, H.W., Cava, R.J., Krajewski, J.J. and Peck Jr., W.J., J. Alloys and Compounds 216, 135 (1994)
- [11] Häse, K., Hough, D., Holzapfel, B. and Schultz, L., Physica B 284-288, 1105 (2000)
- [12] Du Mar, A.C., Rathnayaka, K.D.D., Naugle, D.G. and Canfield, P.C., International Journal of Modern Physics B 12, 3264-3266 (1998)
- [13] Shulga, S.V., Drechsler, S.-L., Fuchs, G., Müller, K.H., Winzer, K., Heinecke, M. and K. Krug, Phys. Rev. Lett. 80, 1730 (1998)
- [14] Singh, D.J., Solid State Commun., 899 (1998)
- [15] Drechsler, S.-L. *et al.*, Physica C 317-318, 117 (1999)